

Mean Field Asymptotics in Statistical Learning.

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Lecture 2: Basics concepts in statistical physics.

Chapter 2
[MM08]

Thermodynamics: heat temperature work.
Macroscopic. 4 laws

Statistical physics: microscopic.
Static dynamics

Most important message: The free energy function is most important quantity to calculate in SP. It can be used to derive most of the properties of physical system.

① Configuration space and Gibbs distribution.

Concept.

Configuration space: Ω

Configuration: $\sigma \in \Omega$

Observable: $f: \Omega \rightarrow \mathbb{R}$

Energy function: $H: \Omega \rightarrow \mathbb{R}$
(Hamiltonian)

Example (Ising model).

$$\Omega = \{\pm 1\}^n \quad (\Omega = \mathbb{R}^n)$$

$$\sigma = \{+1, -1, +1, +1, -1, \dots\}$$

$$f(\sigma) = \sum_{i=1}^n \sigma_i$$

$$H(\sigma) = \sum_{\substack{i,j \in [n] \\ i \neq j}} J_{ij} \sigma_i \sigma_j$$

Ferromagnetic model: $J_{ij} = -1$

Spin glass model: $J_{ij} \stackrel{\text{iid}}{\sim} N(0, \sigma^2)$

★ Gibbs / Boltzmann distribution.

ν_0 is reference measure on Ω .

$$P_\beta(d\sigma) = \frac{1}{Z(\beta)} \exp\{-\beta H(\sigma)\} \nu_0(d\sigma).$$

$$Z(\beta) = \int_{\Omega} \exp\{-\beta H(\sigma)\} \nu_0(d\sigma) \quad \text{Partition function.}$$

β : inverse temperature. $T = 1/\beta$: temperature.

Ensemble average: $f: \Omega \rightarrow \mathbb{R}$. $\langle \cdot \rangle_\beta$

$$\langle f \rangle_\beta \equiv \int_{\Omega} f(\sigma) P_\beta(d\sigma),$$

High temperature limit: $\beta \rightarrow 0$, $T \rightarrow \infty$ $P_\beta \rightarrow \nu_0$ ($\beta \rightarrow 0$)

low temperature limit: $\beta \rightarrow \infty$, $T \rightarrow 0$.

P_β concentrate on $\Omega_0 \equiv \arg\min_{\sigma} H(\sigma)$

$$\lim_{\beta \rightarrow \infty} P_\beta(\Omega_0) = 1.$$

② Thermodynamic potentials. (Special functions of β).

Fix $H = H_n$ P_β

Free energy: $F(\beta) = -\frac{1}{\beta} \log Z(\beta).$

Free entropy: $\Phi(\beta) = \log Z(\beta).$

Internal energy: $U(\beta) = \langle H \rangle_\beta = \int_{\Omega} H(\sigma) P_\beta(d\sigma).$

Canonical entropy: $S(\beta) = -\sum_{\sigma \in \Omega} P_\beta(\sigma) \log P_\beta(\sigma)$

$$a) \quad \Phi'(\beta) = \frac{d}{d\beta} [\log Z(\beta)] = \frac{Z'(\beta)}{Z(\beta)} = -\langle H \rangle_\beta = -U(\beta).$$

$$\begin{aligned} Z'(\beta) &= \frac{d}{d\beta} \int \exp\{-\beta H(\sigma)\} \nu_0(d\sigma) \\ &= \int \exp(-\beta H(\sigma)) \times (-H(\sigma)) \nu_0(d\sigma). \end{aligned}$$

$$b). \Phi''(\beta) = \text{Var}_\beta(H) = [\langle H^2 \rangle_\beta - \langle H \rangle_\beta^2]$$

$$\Phi''(\beta) = \frac{d}{d\beta} \left[\frac{Z'(\beta)}{Z(\beta)} \right] = \frac{Z''(\beta)Z(\beta) - Z'(\beta)^2}{Z(\beta)^2}$$

$$\begin{aligned} Z''(\beta) &= \frac{d}{d\beta} \int \exp(-\beta H(\sigma)) \times (-H(\sigma)) v_0(d\sigma) \\ &= \int \exp(-\beta H(\sigma)) (H(\sigma))^2 v_0(d\sigma). \end{aligned}$$

$$\Phi''(\beta) = \langle H^2 \rangle_\beta - \langle H \rangle_\beta^2 \geq 0.$$

$$\begin{aligned} c) \quad S(\beta) &= - \sum P_\beta(\sigma) \log P_\beta(\sigma) \\ &= - \sum P_\beta(\sigma) [-\beta H(\sigma) - \log Z(\beta)] \\ &= \beta \langle H \rangle_\beta + \log Z(\beta) \\ &= \beta \cdot U(\beta) + \Phi(\beta) \end{aligned}$$

$$F(\beta) = -\frac{1}{\beta} \Phi(\beta) = U(\beta) - \frac{1}{\beta} S(\beta)$$

③ Thermodynamic limit and phase transition.

H_n

$$F_n(\beta), \Phi_n(\beta), U_n(\beta), S_n(\beta).$$

For large n , the potentials are often proportional to n

$$\underline{f(\beta)} = \lim_{n \rightarrow \infty} F_n(\beta) / n \quad \text{free energy density.}$$

$$\phi(\beta) = \lim_{n \rightarrow \infty} \Phi_n(\beta) / n.$$

$$u(\beta) = \lim_{n \rightarrow \infty} U_n(\beta) / n.$$

$$s(\beta) = \lim_{n \rightarrow \infty} S_n(\beta) / n.$$

$f(\beta)$ will be always concave.

$f(\beta)$ is often analytic in some region of β

At some β_c $f(\beta)$ can be non-analytic.

Phase transition at β_c .

④ Ensemble average of an observable.

$$M(\sigma) = \sum_{i=1}^n \sigma_i.$$

Goal: to calculate $\lim_{n \rightarrow \infty} \langle M \rangle_{\beta}^H / n$.

Some oracles: Given any Hamiltonian H_λ we are able

to calculate $F(\beta, \lambda) = -\frac{1}{\beta} \log \int \exp\{-\beta H_\lambda(\sigma)\} \nu_0(d\sigma)$.

Idea: Introduce a perturbed system

$$H_\lambda(\sigma) = H(\sigma) + \lambda M(\sigma).$$

$$P_{\beta, \lambda}(d\sigma) = \frac{1}{Z(\beta, \lambda)} \exp\{-\beta H_\lambda(\sigma)\} \nu_0(d\sigma).$$

$$Z(\beta, \lambda) = \int \exp(\cdot) \nu_0(d\sigma)$$

$$\Phi(\beta, \lambda) = \log Z(\beta, \lambda).$$

$$F(\beta, \lambda) = -\frac{1}{\beta} \log \Phi(\beta, \lambda).$$

$$\langle g \rangle_{\beta, \lambda} = \int_{\Omega} g(\sigma) P_{\beta, \lambda}(d\sigma).$$

★ Prop: $\partial_\lambda F(\beta, \lambda) = \langle M \rangle_{\beta, \lambda}$

$$\partial_\lambda F(\beta, 0) = \langle M \rangle_\beta$$

Rmk: $f(\beta, \lambda) = \lim_{n \rightarrow \infty} F_n(\beta, \lambda) / n$, $m(\beta, \lambda) = \lim_{n \rightarrow \infty} \langle M \rangle_{\beta, \lambda} / n$

$$\partial_\lambda f(\beta, \lambda) = m(\beta, \lambda).$$

⑤ Ensemble variance of an observable.

Lemma: $\partial_\lambda^2 \Phi(\beta, \lambda) = \beta^2 \left[\langle M^2 \rangle_{\beta, \lambda} - \langle M \rangle_{\beta, \lambda}^2 \right]$
 $= \beta^2 \text{Var}_\beta(M)$

