Finite Markov Information-Exchange processes

David Aldous

February 2, 2011

Markov Chains

The next few lectures give a brisk discussion of

- Basics: discrete- and continuous-time.
- Hitting times and mixing times.
- Three standard examples.
- Other examples.

Only occasional math arguments here, **but** when we use some technique later for FMIEs that parallels a technique for MC, we'll recall the MC argument then.

A discrete-time MC (Z(t), t = 0, 1, 2, ...) is specified by its transition matrix $\mathbf{P} = (p_{ij})$. The *t*-step transition probability are

 $\mathbb{P}_i(Z(t) = j)$ are entries of \mathbf{P}^t .

A continuous-time MC (Z(t), $0 \le t < \infty$) is specified by its transition rate matrix $\mathcal{N} = (\nu_{ij})$, where given the off-diagonal entries we set

$$\nu_i := \sum_{j \neq i} \nu_{ij}, \quad \nu_{ii} := -\nu_i.$$

The time-t transition probability are

$$\mathbb{P}_i(Z(t) = j)$$
 are entries of $\exp(\mathcal{N}t)$.

If **irreducible** then there is a unique **stationary distribution** $\pi = (\pi_i)$ and

$$\mathbb{P}_i(Z(t)=j) o \pi_j$$
 as $t o \infty$

holds always in continuous-time. If we can find a distribution $\pi = (\pi_i)$ such that

$$\pi_i
u_{ij} = \pi_j
u_{ji} \,\, \forall j
eq i$$

then π is the stationary distribution and the chain is called **reversible**.

In our FMIE setting we use a symmetric matrix \mathcal{N} and so we have an associated continuous-time MC which is reversible and has **uniform** stationary distribution. We sometimes impose assume **regularity**:

$$\sum_{j\neq i}\nu_{ij}=1 \,\,\forall i. \tag{1}$$

This restriction is loosely analogous to discrete-time random walk on a graph being restricted to a regular graph. Consider the (continuous-time) quantity

$$z_{ij} = \int_{t=0}^{\infty} \left(\mathbb{P}_i(Z(t) = j) - \pi_j \right) dt$$
(2)

analogous to the discrete-time quantity

$$z_{ij} = \sum_{t=0}^{\infty} \left(p_{ij}^{(t)} - \pi_j \right) \tag{3}$$

which can be viewed as a generalized inverse of the singular matrix $I - \mathbf{P}$. The matrix of mean hitting times $\mathbb{E}_i T_j^{\text{hit}}$ can be expressed in terms of the matrix \mathbf{Z} .

Theorem (Mean hitting time formula)

Without assuming reversibility, $\mathbb{E}_i T_j^{hit} = (z_{jj} - z_{ij})/\pi_j$.

See RWG 2.2 for proof and detailed discussion. Here, let me observe three consequences.

Noting $\sum_j z_{ij} = 0 \,\,\forall i$ we get $\sum_j (\mathbb{E}_i T_j^{hit}) \pi_j = \sum_j z_{jj}$ and in particular

Corollary (Random Target Lemma)

 $au_{hit} := \sum_{j} (\mathbb{E}_{i} T_{j}^{hit}) \pi_{j}$ does not depend on i

and so this particular statistic τ_{hit} is the mathematically natural way to summarize the matrix of mean hitting times by a single number. Note however that $\sum_{i} \pi_{i} z_{ij} = 0 \quad \forall j$ and so

$$\mathbb{E}_{\pi} T_j^{\text{hit}} := \sum_i \pi_i \mathbb{E}_i T_j^{\text{hit}} = z_{jj} / \pi_j$$
(4)

which in general **does** depend on *j*.

In the FMIE setting \mathcal{N} is a symmetric matrix, which makes **Z** a symmetric matrix, as well as making π be the uniform distribution. But this does not imply that $(\mathbb{E}_i T_i^{\text{hit}})$ is symmetric; in fact

$$\mathbb{E}_i T_j^{\mathsf{hit}} = \mathbb{E}_j T_i^{\mathsf{hit}} ext{ iff } z_{jj} = z_{ii} ext{ iff } \mathbb{E}_\pi T_j^{\mathsf{hit}} = \mathbb{E}_\pi T_i^{\mathsf{hit}}.$$

A chain is **transitive** if for each pair i_1, i_2 there is a permutation σ of the state space such that $\sigma(i_1) = i_2$ and

$$\nu_{ij} = \nu_{\sigma(i),\sigma(j)} \ \forall i,j.$$

Informally, the chain "looks the same from each state". Transitivity implies $z_{ii} = z_{jj} \forall i, j$ and hence $T_{ij} \stackrel{d}{=} T_{ji} \forall i, j$. [Board: degrees of freedom].

[RWG Chapter 3]

There are two parallel ways to think about the dynamics of the distribution of Z(t). First, in terms of **matrices**. The transition rate matrix \mathcal{N} has eigenvalues

$$0=\lambda_1>-\lambda_2\geq-\lambda_3\geq\ldots\geq-\lambda_n$$

and there is a **spectral representation** (matrix diagonalization)

$$\mathbb{P}_{i}(Z(t) = j) = \pi_{i}^{-1/2} \pi_{j}^{1/2} \sum_{m=1}^{n} \exp(-\lambda_{m} t) u_{im} u_{jm}$$
(5)

for orthonormal \mathbf{U} . In particular, the time-asymptotics for convergence to stationarity are

$$\mathbb{P}_i(Z(t)=j)-\pi_j=c_{ij}e^{-\lambda_2 t}+o(e^{-\lambda_2 t}) \text{ as } t\to\infty. \tag{6}$$

Jargon: λ_2 is the spectral gap, $\tau_{rel} := 1/\lambda_2$ is the relaxation time.

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Knowing λ_2 doesn't tell you anything precise about the finite-time distribution of the MC starting at an arbitrary state, but it does tell you some things about the *stationary* chain. For instance (cf. the extremal characterization later)

$$\max_{f,g} \operatorname{cor}_{\pi}(f(Z(0)),g(Z(t))) = \exp(-\lambda_2 t).$$

Note also

$$\mathbb{P}_i(Z(t) = i) = \pi_i + \sum_{m \ge 2} u_{im}^2 \exp(-\lambda_m t)$$
(7)

so the right side is decreasing with t, and in fact is *completely monotone*.

The second way – which we might call " L^2 theory" or "the Dirichlet formalism" – requires some notational setup. For a "test function" $g : \text{Agents} \to \mathbb{R}$ write

$$\begin{split} \bar{g} &= \sum_{i} \pi_{i} g_{i} \\ ||g||_{2}^{2} &= \sum_{i} \pi_{i} g_{i}^{2} \\ \mathcal{E}(g,g) &= \frac{1}{2} \sum_{i} \sum_{j \neq i} \pi_{i} \nu_{ij} (g_{j} - g_{i})^{2} \quad \text{(the Dirichlet form).} \end{split}$$

When $\bar{g} = 0$ then $||g||_2$ measures "global" variability of g whereas $\mathcal{E}(g,g)$ measures "local" variability relative to the underlying geometry. [Discussion on board]

For a (signed) measure θ we define $||\theta||_{2(m)} = ||f||_2$ for the density $f_i = \theta_i / \pi_i$, and then for a PM μ we have

$$|\mu - \pi||_{2(m)}^2 = -1 + \sum_i \mu_i^2 / \pi_i.$$

This is " L^2 distance" for probability measures.

Why is this viewpoint useful?

The basic evolution equation, for the time-t distribution $\rho(t) = (\rho_j(t))$ from an arbitrary start, is

$$\frac{d}{dt}\rho_j(t) = \sum_j \nu_{ij}\rho_i(t) \tag{8}$$

from which we previously obtained

 $\mathbb{P}_i(Z(t) = j)$ are entries of $\exp(\mathcal{N}t)$.

But a little algebra, directly from (8), gives

Lemma (Global convergence equation)

$$rac{d}{dt} \|
ho(t) - \pi \|_{2(m)}^2 = -2 \mathcal{E}(f(t), f(t)); \qquad f_j(t) =
ho_j(t)/\pi_j.$$

Because $\mathcal{E} \geq 0$ this gives a certain "monotonicity" in global convergence; cf. monotonicity of $\mathbb{P}_i(Z(t) = i)$.

Reformulating the classical *Rayleigh–Ritz* extremal characterization of eigenvalues:

Theorem (Extremal characterization of relaxation time)

 $au_{rel} = \sup\{||g||_2^2 / \mathcal{E}(g,g): \ \bar{g} = 0\}.$

So we can get lower bounds on τ_{rel} by plugging in a test function g chosen heuristically. Much of (algorithm-related) uses of finite MCs involves getting reasonable upper bounds on τ_{rel} and τ_{mix} below. The extremal characterization doesn't help directly but is the starting point for other methodologies.

Combining the extremal characterization with the global convergence equation leads easily to [calculation on board]

Lemma (L^2 contraction lemma)

The time-t distributions $\rho(t)$ of a reversible MC satisfy

$$\|
ho(t)-\pi\|_{2(m)}\leq e^{-t/ au_2}\|
ho(0)-\pi\|_{2(m)}.$$

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Hitting times and mixing times are distinct aspects of a MC; but here's a minor connection. Write T_A^{hit} for the hitting time on a subset $A \subset \text{Agents}$.

Proposition

For a subset A of a continuous-time chain,

$$\sup_{t} |\mathbb{P}_{\pi}(T_{A} > t) - \exp(-t/\mathbb{E}_{\pi}T_{A})| \leq \tau_{rel}/\mathbb{E}_{\pi}T_{A}.$$

In words: for the hitting time distribution to be approximately Exponential it is sufficient that the mean hitting time be large compared to the relaxation time $\tau_{\rm rel}$.

Theory project. Give a bound on the dependence between initial state X(0) and T_A , for instance

$$\max_{f,g} \operatorname{cor}(f(X(0)), h(T_A)) \leq \psi(\tau_{\mathsf{rel}} / \mathbb{E}_{\pi} T_A).$$

For another connection, recall the *Random Target Lemma* said that $\tau_{\text{hit}} := \sum_{j} (\mathbb{E}_{i} T_{j}^{\text{hit}}) \pi_{j}$ does not depend on *i*. It turns out that τ_{hit} has a simple expression in terms of the eigenvalues:

$$au_{\text{hit}} = \sum_{i=1}^{n} 1/\lambda_i$$
 the eigentime identity. (9)

Variation distance (or **total variation**) between a PM μ and the stationary distribution π is defined as

$$||\mu - \pi||_{VD} := \frac{1}{2} \sum_{i} |\mu_i - \pi_i|.$$

This is essentially " L^1 distance". Note that, like " L^2 distance", it ignores the geometry.

For a continuous-time MC, define (variation distance) mixing time τ_{mix} to be the smallest time *t* for which

$$\max ||\mathbb{P}_i(Z(t)\in \cdot)-\pi(\cdot)||_{\mathit{VD}}\leq 1/(2e).$$

The choice of constant on the right must be < 1/2 but is otherwise rather arbitrary; the particular choice 1/(2e) ensures

$$au_{\mathsf{rel}} \leq au_{\mathsf{mix}}.$$

Variation distance and τ_{mix} are central to many theoretical algorithmic uses of MCs – see Montenegro-Tetali (2006) and the monographs.

The general version of the "bottleneck parameters" earlier are defined in terms of stationary flow rates

$$Q(A, A^c) := \sum_{i \in A, j \in A^c} \pi_i \nu_{ij}.$$

In particular, define the Cheeger time constant by

$$au_{\mathsf{cond}} := \sup_{A} rac{\pi(A)(1-\pi(A))}{Q(A,A^c)}.$$

[Discussion on board: up to factors of 2 this is 1/conductance; *n*-cycle case].

There is a (not easy)

Theorem (Cheeger's inequality)

For any continuous-time reversible MC,

$$\tau_{rel} \le 8\tau_{cond}^2 \max_i \nu_i.$$

So in the FMIE context with the regularity assumption (1) we have

$$\phi(m) = \min\{\nu(A, A^c) : |A| = m\}, \quad 1 \le m \le n-1$$
$$\tau_{\text{cond}} := \sup_{m} \frac{\frac{m}{n} \frac{n-m}{n}}{\phi(m)}.$$

and Cheeger's inequality becomes

$$au_{\mathrm{rel}} \leq 8 au_{\mathrm{cond}}^2.$$

More sophisticated results can be found in the survey by Montenegro-Tetali (2006).

Helpful intuition is that a sequence

$$\hat{Z}_1, \hat{Z}_2, \ldots$$

obtained as either a stationary MC sampled at multiples of $\tau_{\rm rel}$

$$Z(\tau_{\mathsf{rel}}), Z(2\tau_{\mathsf{rel}}), \ldots$$

or an arbitrary-start MC sampled at multiples of $\tau_{\rm mix}$

$$Z(\tau_{\min}), Z(2\tau_{\min}), \ldots$$

behaves similarly to an IID sequence as far as quantitative versions of limit theorems are concerned. See e.g. León-Perron (2004) for a large deviation inequality for occupation times.

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Another well-studied MC topic is the cover time

$$C := \max_{j} T_{j}^{hit} = time until every state visited.$$

There is a "non-clever" bound in term of the parameter $\tau^* := \max_{i,j} \mathbb{E}_i T_j^{\text{hit}}$, because inductively

$$\mathbb{P}_i(T_j^{\mathsf{hit}} > 2m\tau^*) \le 2^{-m}, \ m = 1, 2, 3, \dots$$

and it quickly follows that

$$\max_i \mathbb{E}_i C \leq (1+o(1))\tau^* \log n.$$

And a "clever" argument called Matthews' method sharpens this to

$$\max_{i} \mathbb{E}_{i} C \leq \tau^* \sum_{i=1}^{n-1} 1/i.$$

Recent deep results of Ding-Lee-Peres (2010) give very sharp general estimates of $\mathbb{E}C$.

Reversible Markov chains: standard examples

For any discrete-time MC with transition probabilities p_{ij} there is a corresponding continuous-time MC with transition rates $\nu_{ij} = p_{ij}$. In particular, discrete-time RW on a *d*-regular undirected graph is the MC with transition probabilities

$$p_{ij} = d^{-1}$$
 for edges (i, j)

and there is a corresponding continuous-time RW.

For a continuous-time MC, in the case where $\nu_i := \sum_{j \neq i} \nu_{ij}$ is constant in *i*, it is natural to standardize the time unit so that $\nu_i \equiv 1$.

[board: comments re 2 different continuous-times RWs on graphs with highly varying degrees – needs watching in all FMIE contexts]

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Continuous-time RW on the complete *n*-vertex graph.

$$u_{ij}=1/(n-1), \quad j\neq i.$$

The basics are easy - no surprise!

$$\mathbb{E}_{i} T_{j}^{\text{hit}} = n - 1; \quad T_{j}^{\text{hit}} \stackrel{d}{=} \text{Exponential}(1/(n - 1)).$$
$$\mathbb{P}_{i}(Z(t) = i) = \frac{1}{n} + (1 - \frac{1}{n}) \exp\left(-\frac{n}{n-1}t\right).$$
$$\tau_{\text{rel}} = \frac{n-1}{n}.$$

Because here $\tau_{\rm rel} \approx 1$, for other geometries we can think of $\tau_{\rm rel}$ as relaxation time relative to the complete graph case.

Continuous-time RW on the *d*-dimensional lattice/cube/torus.

First consider the infinite lattice \mathbb{Z}^d . Discrete-time RW on \mathbb{Z}^d is a well-studied classical object. The continuous-time RW $Z^{(d)}(t)$ is nicer in that the co-ordinate processes are independent slowed-down 1-dimensional RWs; for the origin **0** and $\mathbf{x} = (x_1, \ldots, x_d)$

$$\mathbb{P}_{\mathbf{0}}(Z^{(d)}(t) = \mathbf{x}) = \prod_{i=1}^{d} \mathbb{P}_{\mathbf{0}}(Z^{(1)}(t/d) = x_i).$$

Five facts you should know about RW on \mathbb{Z}^d .

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(CLT): The distribution of $Z^{(d)}(t)$ for large t is approximately multivariate Normal; marginals are Normal(0, t/d). (Local density) $\mathbb{P}_0(Z^{(1)}(t) = 0) \sim (2\pi)^{-1/2}t^{-1/2}$ and so

$$\mathbb{P}_{\mathbf{0}}(Z^{(d)}(t) = \mathbf{0}) \sim (2\pi t/d)^{-d/2}.$$

(Recurrence/transience) In d = 1, 2 RW is **recurrent**: each vertex is visited infinitely often. In $d \ge 3$ RW is **transient**: the chance state **x** is ever visited $\rightarrow 0$ as $|\mathbf{x}| \rightarrow \infty$.

(Fair game: winner and mean duration). In d = 1, for -a < 0 < b

$$\mathbb{P}_0(extsf{T}_b < extsf{T}_{-a}) = a/(a+b); \quad \mathbb{E}_0\min(extsf{T}_{-a}, extsf{T}_b) = ab$$

Continuous-time RW on the *d*-dimensional torus.

One natural geometry is the 2-dimensional discrete square $[0, m-1]^2$ as a subgraph of \mathbb{Z}^2 . It is mathematically nicer to eliminate the boundary by imposing "periodic boundary conditions", that is to use the 2-dimensional discrete torus, which is vertex-transitive. In general dimension $d \ge 1$ this becomes the *d*-dimensional (discrete) torus, denoted \mathbb{Z}_m^d .

Warning; we study $m \to \infty$ asymptotics for fixed *d*. To compare with other models, remember $n = m^d$.

We quote some basic facts. Consider d = 1, so \mathbb{Z}_m is the *m*-cycle. The eigenvalues are

$$\cos(2\pi j/m), \ 0 \leq j \leq m-1$$

and the relaxation time is

$$au_{
m rel} = rac{1}{1-\cos(2\pi/m)} \sim rac{m^2}{2\pi^2}.$$

For $d \ge 2$ we retain the nice property that the co-ordinate processes are independent slowed-down versions of the RW on the *m*-cycle; so for fixed *m* the *d*-dimensional and 1-dimensional RWs are again related by

$$\mathbb{P}_{\mathbf{0}}(Z^{(d)}(t) = \mathbf{x}) = \prod_{i=1}^{d} \mathbb{P}_{\mathbf{0}}(Z^{(1)}(t/d) = x_i).$$

From this we see that the eigenvalues on \mathbb{Z}_m^d are

$$\lambda_{(k_1...k_d)} = \frac{1}{d} \sum_{u=1}^d (1 - \cos(2\pi k_u/m)), \ 0 \le k_u \le m - 1.$$

In particular, the relaxation time satisfies

$$\pi_2 \sim rac{dm^2}{2\pi^2} = rac{dn^{2/d}}{2\pi^2}.$$

We can also use the *eigentime identity* to compute the mean hitting time parameter

$$au_{\mathsf{hit}} = \sum_{k_1} \cdots \sum_{k_d} 1/\lambda_{(k_1,...,k_d)}$$

(the sum excluding $(0, \ldots, 0)$),

and hence

$$au_{\rm hit} \sim m^d R_d$$
 (10)

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where

$$R_{d} \equiv \int_{0}^{1} \cdots \int_{0}^{1} \frac{1}{\frac{1}{d} \sum_{u=1}^{d} (1 - \cos(2\pi x_{u}))} dx_{1} \cdots dx_{d}$$
(11)

provided the integral converges. In fact by the recurrence/transience properties of RW on the whole integer lattice we must have $R_d < \infty$ for $d \geq 3$ only. For d = 1 we must have $\tau_{hit} = \Theta(m^2)$, and the case d = 2 is best understood via a later argument.

Random graphs with prescribed degree distributions

[on board: only key points here]

- Maybe 500 papers since 2000 on such random graph models.
- Configuration model: basic properties and local Galton-Watson approximation.
- Continuous-time vs discrete-time RW
- Mean hitting times via tree recursions. In particular, on random *r*-regular graph $\tau_{hit} \sim \frac{r-1}{r-2}n$

Outline. Specify (d_i) , Can define models \mathcal{G}_n of *n*-vertex graph, interpretable as being "random" subject to the following constraint. Write D_n for degree of a uniform random vertex of \mathcal{G}_n , then

$$D_n \stackrel{d}{\rightarrow} D$$
 where $\mathbb{P}(D=i) = d_i$.

Such models have the following "local GWBP approximation". The structure of \mathcal{G}_n within some fixed graph-distance r from a uniform random vertex U_n converges in distribution, as $n \to \infty$, to the random tree comprising generations 0 to r of the following modified Galton-Watson BP. The root has offspring distribution D; in subsequent generation the offspring distribution is the size-biased distribution D^* where $\mathbb{P}(D^* = i) = (i+1)d_{i+1}/\mathbb{E}D$.

Assuming $d_0 = d_1 = 0$, the GWBP is an infinite tree (non-extinction). Assuming $\mathbb{E}D^{2+\varepsilon} < \infty$ then $\mathbb{E}(D^*)^{1+\varepsilon} < \infty$ and the Kesten-Stigum theorem says that the size Y_r of generation r grows at a particular rate: $Y_r/(\mathbb{E}D^*)^r \to W$ a.s. and L^1 .

The results above suggest **heuristics** for the structure of \mathcal{G}_n and the behavior of RW and other FMIE processes on \mathcal{G}_n .

Let us record the following **local transience principle.** For a large finite-state MC whose behavior near a state *i* can be approximated be a transient infinite-state chain,

$$\mathbb{E}_{\pi} T_i^{\mathsf{hit}} pprox R_i / \pi_i$$

where R_i is defined in terms of the approximating infinite-state chain as $\int_0^\infty p_{ii}(t) dt = \frac{1}{\nu_i q_i}$, where q_i is the chance the infinite-state chain started at *i* will never return to *i*.

The approximation comes from (4) via a "interchange of limits" procedure which requires ad hoc justification.

In the case of simple RW on the $d \ge 3$ -dimensional torus, this identifies the constant R_d at (11) as $R_d = 1/q_d$ where q_d is the chance that RW never returns to the origin. So (11) provides a formula for q_d .

In the "random graphs with prescribed degree distribution" model, this argument shows (heuristics) that $\mathbb{E}_{\pi} T_i^{\text{hit}} = \Theta(n)$.

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Other geometries

The *d*-dimensional hypercube $\{0,1\}^d$ is often used as the simplest non-trivial example of a geometry on which the RW is rapidly mixing. In particular it illustrates the cut-off window for variation distance mixing [xxx explain on board]. But it seems not so natural for the applications we have in mind.

Small world graphs, which start with the *d*-dimensional lattice and add random long edges (v, w) with probabilities $\propto ||v - w||^{-\gamma}$, are interesting but hard to study analytically.

Proximity graphs, described next, are also interesting but hard to study analytically.

Given points (x_i) in the plane in general position, create edges according to a deterministic rule such as

create an edge (x_i, x_j) iff the disc $A(x_i, x_j)$ with diameter-line (x_i, x_j) does not contain any third point of (x_i) .



Replacing the disc with a one-parameter family of other shapes, and applying this construction to random (Poisson) points, gives a family of **random proximity graphs** which (unlike the more familiar **random geometric graphs**) are always connected.

Simulation project. Study RW and other FMIE processes on these graphs.