# The Zero Set and Arcsine Laws of Brownian Motion 

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In this lecture, we will consider the properties of the zero set of Brownian motion and introduce two arcsine laws. Throughout this lecture, let $\left(B_{t}, t \geq 0\right)$ be the Brownian motion starting from 0. $P_{x}$ is the distribution of Brownian motion starting from $x$.

First of all, we define the zero set of $\left(B_{t}, t \geq 0\right)$ as $\mathcal{Z}(\omega):=\left\{t: B_{t}(\omega)=0\right\}$. Since $B$ has continuous paths, $\mathcal{Z}(\omega)$ is closed subset of $[0, \infty)$, which depends on $\omega$ through the path of $B$. Intuitively, $\mathcal{Z}(\omega)$ is a random closed subset of $[0, \infty)$. Exercise: make this rigorous by putting an appropriate $\sigma$-field on on the set of all closed subsets of $[0, \infty)$.

Now let $|\mathcal{Z}(\omega)|$ denote the Lebesgue measure of $\mathcal{Z}(\omega)$. The first result is
Theorem 18.1. $|\mathcal{Z}(\omega)|=0$ almost surely.

Proof. For any $t \neq 0$, since $B_{t} \sim N(x, t)$ under $P_{x}$, we have

$$
P_{x}(t \in \mathcal{Z})=P_{x}\left(B_{t}=0\right)=0
$$

i.e.

$$
\int 1_{t \in \mathcal{Z}(\omega)} d P_{x}(\omega) \quad \forall t \neq 0
$$

By Fubini's theorem, we get the following relation,

$$
E(|\mathcal{Z}|)=\iint_{0}^{\infty} 1_{t \in \mathcal{Z}(\omega)} d t d P_{x}(\omega)=\int_{0}^{\infty} \int 1_{t \in \mathcal{Z}(\omega)} d P_{x}(\omega) d t=0
$$

which implies $|\mathcal{Z}|=0$ a.s.

Another property of $\mathcal{Z}$ is
Theorem 18.2. With probability one, $\mathcal{Z}(\omega)$ has no isolated points.

Proof. Let $R_{t}=\inf \left\{u>t: B_{u}=0\right\}, T_{0}=\inf \left\{u>0: B_{u}=0\right\}$. The recurrence implies $P_{x}\left(R_{t}<\infty\right)=1$, so by the strong Markov property, we know

$$
P_{x}\left(T_{0} \circ \theta_{R_{t}}>0 \mid \mathcal{F}_{R_{t}}\right)=P_{0}\left(T_{0}>0\right)=0
$$

Take expectation again we get

$$
P_{x}\left(T_{0} \circ \theta_{R_{t}}>0 \text { for some } t \in \mathbb{Q}\right)=0
$$

So if a point $u \in \mathcal{Z}(\omega)$ is isolated from the left, i.e. $u=R_{t}$ for some rational $t$, calculation above shows that it is an accumulated point from the right. So $\mathcal{Z}(\omega)$ has no isolated points a.s.

Now we come to see the arcsine laws for Brownian motion. There are at least three different arcsine laws. Here we will introduce two of them.

Let $T=\arg \max _{0<t<1} B_{t}$, notice $T$ is well defined because all the local maximum of Brownian motion are distinct. Our first arcsine law is

Theorem 18.3. For any $t \in[0,1], P_{0}(T \leq t)=\frac{2}{\pi} \arcsin (\sqrt{t})$.

Proof. For any $t \in[0,1]$, let $X_{r}=B_{t-r}-B_{t}, Y_{s}=B_{t+s}-B_{t}$, then $\left(X_{r}, 0 \leq r \leq t\right)$ is a Brownian motion starting from 0 and is $\mathcal{F}_{t}$ measurable; $\left(Y_{s}, 0 \leq s \leq 1-t\right)$ is also a Browian motion starting from 0 and is independent of $\mathcal{F}_{t}$. With this setup, we have

$$
\begin{aligned}
P_{0}(T \leq t) & =P_{0}\left(\max _{[0, t]} B_{u}>\max _{[t, 1]} B_{u}\right) \\
& =P_{0}\left(\max _{[0, t]}\left(B_{u}-B_{t}\right)>\max _{[t, 1]}\left(B_{u}-B_{t}\right)\right) \\
& =P_{0}\left(\max _{[0, t]}\left(B_{t-r}-B_{t}\right)>\max _{[0,1-t]}\left(B_{t+s}-B_{t}\right)\right) \\
& =P_{0}\left(\max _{[0, t]} X_{r}>\max _{[0,1-t]} Y_{s}\right)
\end{aligned}
$$

By previous work, we know $\max _{[0, t]} X_{r} \stackrel{d}{=}\left|X_{t}\right|, \max _{[0,1-t]} Y_{s} \stackrel{d}{=}\left|Y_{1-t}\right|$, and they are independent. So if we let $Z_{1}, Z_{2}$ are i.i.d. $\sim N(0,1)$, and $\theta$ is uniformly distributed on $[0,2 \pi)$, then

$$
\begin{aligned}
P_{0}(T \leq t) & =P_{0}\left(\left|X_{t}\right|>\left|Y_{1-t}\right|\right) \\
& =P\left(\sqrt{t}\left|Z_{1}\right|>\sqrt{1-t}\left|Z_{2}\right|\right) \\
& =P\left(\frac{\left|Z_{2}\right|}{\sqrt{Z_{1}^{2}+Z_{2}^{2}}}<t\right) \\
& =P(|\sin \theta|<\sqrt{t}) \\
& =\frac{2}{\pi} \arcsin (\sqrt{t})
\end{aligned}
$$

A little more calculation will show us the last zero point of Brownian motion on $[0,1]$ has the same distribution as T above. Let $L=\sup \left\{t \leq 1: B_{t}=0\right\}$, and for $a>0, T_{a}=\inf \left\{t: B_{t}=a\right\}$, by previous work, for any $t>0$,

$$
\begin{aligned}
P_{0}\left(T_{a} \leq t\right) & =2 P_{0}\left(B_{t} \geq a\right)=2 \int_{a}^{\infty}(2 \pi t)^{-1 / 2} \exp \left(-x^{2} / 2 t\right) d x \\
& =2 \int_{t}^{0}(2 \pi t)^{-1 / 2} \exp \left(-a^{2} / 2 s\right)\left(-t^{1 / 2} a / 2 s^{3 / 2}\right) d s \\
& =\int_{0}^{t}\left(2 \pi s^{3}\right)^{-1 / 2} a \exp \left(-a^{2} / 2 s\right) d s
\end{aligned}
$$

Use this formula, we have
Theorem 18.4. For any $s \in[0,1], P_{0}(L \leq s)=\frac{2}{\pi} \arcsin (\sqrt{s})$.

Proof. By Markov property of Brownian motion,

$$
\begin{aligned}
P_{0}(L \leq s) & =\int_{-\infty}^{\infty} p_{s}(0, x) P_{x}\left(T_{0}>1-s\right) d x \\
& =2 \int_{0}^{\infty} p_{s}(0, x) P_{0}\left(T_{x}>1-s\right) d x \\
& =2 \int_{0}^{\infty}(2 \pi s)^{-1 / 2} \exp \left(-x^{2} / 2 s\right) \int_{1-s}^{\infty}\left(2 \pi r^{3}\right)^{-1 / 2} x \exp \left(-x^{2} / 2 r\right) d r d x \\
& =\frac{1}{\pi} \int_{1-s}^{\infty}\left(s r^{3}\right)^{-1 / 2} \int_{0}^{\infty} x \exp \left(-x^{2}(r+s) / 2 r s\right) d x d r \\
& =\frac{1}{\pi} \int_{1-s}^{\infty}\left(s r^{3}\right)^{-1 / 2} r s /(r+s) d r \\
& =\frac{1}{\pi} \int_{0}^{s}(t(1-t))^{-1 / 2} d t \quad(\text { let } t=s /(r+s)) \\
& =\frac{2}{\pi} \arcsin (\sqrt{s})
\end{aligned}
$$

R. Durrett. (1996). Probability: theory and examples.(2nd Edition) Duxbury Press.

