

Lecture 29

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When $h = 0$, a phase transition occurs at $\beta = 1$. We saw that in the high temperature phase, $R_{12} = O(N^{-1/2})$. Parisi conjectures that at $\beta = 1$ and $h = 0$, R_{12} is of order $N^{-1/3}$. Guerra proved that at $\beta = 1$ and $h = 0$, $\mathbf{E}\langle R_{12}^2 \rangle \leq C/\sqrt{N}$ for some constant C , ie, $R_{12} = O(N^{-1/4})$ at most. Talagrand has another proof in his book, but it is complicated. Nothing better is known.

Using Stein's method, we will show that $\mathbf{E}\langle |R_{12}|^3 \rangle \geq C/N$ for some $C > 0$.

Proof: Let $\psi(x)$ be the probability density

$$\frac{\cosh(x)e^{-x^2/2}}{\sqrt{2\pi e}}$$

(i.e. the symmetric mixture of $N(1, 1)$ and $N(-1, 1)$). Hopefully, this is the distribution of the local field as $N \rightarrow \infty$.

For any bounded, measurable f , let $Mf = \int_{-\infty}^{\infty} f(x)\psi(x) dx$. Define an operator U as

$$Uf(x) = \frac{e^{x^2/2}}{\cosh x} \int_{-\infty}^{\infty} \cosh(t)e^{-t^2/2}(f(t) - Mf) dt$$

and let $Tf(x) = f'(x) - (x - \tanh(x))f(x)$. Verify that $TUf = f - Mf$, so U is the inversion of the Stein operator.

Lemma 1 $\|Uf\|_{\infty} \leq C\|f\|_{\infty}$ and $\|(Uf)'\|_{\infty} \leq C\|f\|_{\infty}$ for some universal constant C .

Lemma 2 (Expansion lemma) Fix a bounded, measurable $f : \mathbb{R} \rightarrow \mathbb{R}$, let b_1, \dots, b_m be arbitrary functions of σ . Assume that b_1 does not depend on σ_1 . Then

$$\begin{aligned} \mathbf{E}(\langle f(l_1)b_1 \rangle \langle b_2 \rangle \cdots \langle b_m \rangle) &= (Mf)\mathbf{E}(\langle b_1 \rangle \cdots \langle b_m \rangle) \\ &\quad - \sum_{r=2}^m \frac{1}{N} \sum_{j=2}^N \mathbf{E}(\langle \sigma_j Uf(l_1)b_1 \rangle \langle b_2 \rangle \cdots \langle b_r \sigma_1 \sigma_j \rangle \langle b_{r+1} \rangle \cdots \langle b_m \rangle) \\ &\quad + \frac{m}{N} \sum_{j=2}^N \mathbf{E}(\langle \sigma_j Uf(l_1)b_1 \rangle \langle b_2 \rangle \cdots \langle b_m \rangle) + \frac{1}{N} \mathbf{E}(\langle (Uf)'(l_1)b_1 \rangle \langle b_2 \rangle \cdots \langle b_m \rangle) \\ &\quad + \frac{1}{N} \mathbf{E}(\langle \sigma_j Uf(l_1)b_1 \sigma_1 \rangle \langle b_2 \rangle \cdots \langle b_m \rangle) \end{aligned}$$

Now take $f(x) = \tanh(x)$. Then $\langle \sigma_1 \sigma_2 \rangle = \langle f(l_1) \sigma_2 \rangle$ (since $f(l_1)$ is the conditional expectation of σ_1 given the rest of the spins). Thus $\mathbf{E}\langle \sigma_1 \sigma_2 \rangle^2 = \mathbf{E}\langle f(l_1) \sigma_2 \rangle \langle \sigma_1 \sigma_2 \rangle$.

Let $b_1 = \sigma_2$, $b_2 = \sigma_1 \sigma_2$. Since ψ is a symmetric density and \tanh is odd, $Mf = 0$. Let $h = Uf$. Applying the expansion lemma,

$$\begin{aligned} \mathbf{E}\langle \sigma_1 \sigma_2 \rangle^2 &= 0 - \frac{1}{N} \sum_{j=2}^N \mathbf{E}\langle \sigma_1 h(l_1) \sigma_2 \rangle \langle \sigma_2 \sigma_j \rangle \\ &\quad + \frac{2}{N} \sum_{j=2}^N \mathbf{E}\langle \sigma_j h(l_1) \sigma_2 \rangle \langle \sigma_2 \sigma_j \rangle \langle \sigma_1 \sigma_j \rangle \\ &\quad + \frac{1}{N} \mathbf{E}\langle h'(l_1) \sigma_2 \rangle \langle \sigma_1 \sigma_2 \rangle \\ &\quad + \frac{1}{N} \mathbf{E}\langle h(l_1) \sigma_1 \sigma_2 \rangle \langle \sigma_1 \sigma_2 \rangle \end{aligned}$$

To simplify the above, let Rem denote any term that is bounded by $CN^{-1}\sqrt{\mathbf{E}\langle R_{12} \rangle^2}$ for some constant C .

Since h' is bounded, so

$$\begin{aligned} \frac{1}{N} \mathbf{E}\langle h'(l_1) \sigma_2 \rangle \langle \sigma_1 \sigma_2 \rangle &\leq \frac{C}{N} \mathbf{E}|\langle \sigma_1 \sigma_2 \rangle| \leq \frac{C}{N} \sqrt{\mathbf{E}\langle \sigma_1 \sigma_2 \rangle^2} \\ &\leq \frac{C'}{N} \sqrt{\mathbf{E} \frac{1}{N^2} \sum_{i,j} \langle \sigma_i \sigma_j \rangle^2} = \frac{C'}{N} \sqrt{\mathbf{E}\langle R_{12} \rangle^2}. \end{aligned}$$

So the fourth term is Rem . The last term is also Rem . Any single term in the sum in the third term is also Rem . In the 2nd term, for $j = 2$, we get $\frac{1}{N} \mathbf{E}\langle h(l_1) \rangle$. All other terms are Rem . So

$$E\langle \sigma_1 \sigma_2 \rangle^2 = -\frac{1}{N} \mathbf{E}\langle h(l_1) \rangle - \mathbf{E}\langle \langle h(l_1) \sigma_2 \sigma_3 \rangle \langle \sigma_2 \sigma_3 \rangle \rangle - 2\mathbf{E}\langle \langle h(l_1) \sigma_2 \sigma_3 \rangle \langle \sigma_1 \sigma_2 \rangle \langle \sigma_1 \sigma_3 \rangle \rangle - Rem.$$

This is the first expansion. To complete the proof, we apply the expansion lemma to each of the above terms. It will be enough to show:

$$-\frac{\mathbf{E}\langle h(l_1) \rangle}{N} = \frac{1}{N} + Rem,$$

$$-\mathbf{E}\langle \langle h(l_1) \sigma_2 \sigma_3 \rangle \langle \sigma_2 \sigma_3 \rangle \rangle = E\langle \sigma_1 \sigma_2 \rangle^2 + T_1,$$

where $|T_1| \leq C\mathbf{E}\langle |R_{12}|^3 \rangle$, and

$$|\mathbf{E}\langle \langle h(l_1) \sigma_2 \sigma_3 \rangle \langle \sigma_1 \sigma_2 \rangle \langle \sigma_1 \sigma_3 \rangle \rangle| \leq C\mathbf{E}\langle |R_{12}|^3 \rangle + Rem.$$

If we can show the above, then

$$\mathbf{E}\langle\sigma_1\sigma_2\rangle^2 = \frac{1}{N} + \mathbf{E}\langle\sigma_1\sigma_2\rangle^2 + T_1 + T_2$$

where T_1, T_2 are bounded by $C\mathbf{E}\langle|R_{12}|^3\rangle + Rem$. So

$$\frac{1}{N} \leq C\mathbf{E}\langle|R_{12}|^3\rangle + Rem \leq C\mathbf{E}\langle|R_{12}|^3\rangle + \frac{C'\sqrt{\mathbf{E}\langle R_{12}^2\rangle}}{N}$$

Now $\sqrt{\mathbf{E}\langle R_{12}^2\rangle} \leq (\mathbf{E}\langle|R_{12}|^3\rangle)^{1/3}$. Suppose $\mathbf{E}\langle|R_{12}|^3\rangle \leq \frac{1}{2CN}$. Then we get $1/N \leq 1/2N + (C'/N)(1/2CN)^{1/3}$, a contradiction when N is large enough. So for N large enough, $\mathbf{E}\langle|R_{12}|^3\rangle \geq \frac{1}{2CN}$.

To prove the above three statements:

Apply the approximation lemma, and only the first terms will matter. Let $w = Uh$ (so we invert the Stein operator again). Verify that $Mh = -1$. By the expansion lemma,

$$\begin{aligned} \mathbf{E}\langle h(l_1) \rangle &= -1 + \frac{1}{N} \sum_{j=2}^N \mathbf{E}\langle \langle \sigma_j w(l_1) \rangle \langle \sigma_1 \sigma_j \rangle \rangle + \frac{\mathbf{E}\langle w'(l_1) \rangle + \mathbf{E}\langle \sigma_1 \sigma_j \rangle}{N} \\ &= -1 + Rem + O(1/N). \end{aligned}$$

Using the expansion lemma on the second term, $\mathbf{E}\langle \langle h(l_1) \sigma_2 \sigma_3 \rangle \langle \sigma_2 \sigma_3 \rangle \rangle = (Mh)\mathbf{E}\langle \sigma_2 \sigma_3 \rangle^2 = -\mathbf{E}\langle \sigma_2 \sigma_3 \rangle^2 = -\mathbf{E}\langle \sigma_1 \sigma_2 \rangle^2$ with some remainder terms. The third term can also be bounded through the expansion lemma.

Lastly, a sketch of the proof of the expansion lemma:

We want to find $\mathbf{E}\langle f(l_1) b_1 \rangle \cdots \langle b_m \rangle$. Let $h = Uf$ so that $h'(x) - (x - \tanh(x))h(x) = f(x) - Mf$. Replace $f(l_1) - Mf$ by $h'(l_1) - l_1 h(l_1) + \tanh(l_1) h(l_1)$ and apply integration by parts on the terms arising from $l_1 h(l_1)$.

□

Exercise 3 Get an upper bound for $\mathbf{E}\langle|R_{12}|^3\rangle$.

Exercise 4 Evaluate

$$\lim_{N \rightarrow \infty} \mathbf{E}\langle N \langle R_{12} R_{23} R_{31} \rangle \rangle$$

or, alternatively,

$$\lim_{N \rightarrow \infty} \mathbf{E}\langle N \langle \sigma_1 \sigma_2 \rangle \langle \sigma_2 \sigma_3 \rangle \langle \sigma_3 \sigma_1 \rangle \rangle$$

You can use Guerra's result.